

$$\begin{aligned}
 \Gamma(z) &:= I_K - \sum_{i=1}^{p-1} \Gamma_i z^i, \\
 B_*(z) &:= Q[\Gamma(z)\bar{\beta}(1-z) - \alpha z : \Gamma(z)\beta_{\perp}], \\
 B(z) &= I_K - \sum_{i=1}^p B_i z^i := Q^{-1}B_*(z)Q,
 \end{aligned} \tag{6.3.13}$$

and

$$\Theta(z) := B(z)^{-1} = \sum_{j=0}^{\infty} \Theta_j z^j.$$

Notice that  $B(0) = Q^{-1}B_*(0)Q = [\bar{\beta} : \beta_{\perp}]Q = I_K$ . Hence,  $B(z)$  has the representation  $I_K - \sum_{i=1}^p B_i z^i$  stated in (6.3.13). Moreover, the matrix operator  $\Theta(z)$  can be decomposed as

$$\Theta(z) = \Theta(1) + (1-z)\Theta^*(z),$$

where expressions for the  $\Theta_j^*$ 's can be found by comparing coefficients in  $\Theta(z) = \sum_{j=0}^{\infty} \Theta_j z^j$  and

$$\begin{aligned}
 \Theta(1) + (1-z)\Theta^*(z) &= \Theta(1) + \sum_{j=0}^{\infty} \Theta_j^* z^j (1-z) \\
 &= (\Theta(1) + \Theta_0^*) + \sum_{j=1}^{\infty} (\Theta_j^* - \Theta_{j-1}^*) z^j.
 \end{aligned}$$

Hence,

$$\Theta_0 = \Theta(1) + \Theta_0^*$$

and

$$\Theta_i = \Theta_i^* - \Theta_{i-1}^*, \quad i = 1, 2, \dots$$

Using the last expression, we get by successive substitution,

$$\begin{aligned}
 \Theta_i^* &= \Theta_i + \Theta_{i-1}^* = \sum_{j=1}^i \Theta_{i-j} + \Theta_0^* \\
 &= \sum_{j=1}^i \Theta_{i-j} + \Theta_0 - \Theta(1) = - \sum_{j=i+1}^{\infty} \Theta_j, \quad i = 1, 2, \dots
 \end{aligned} \tag{6.3.14}$$

From these quantities the operator  $\Xi^*(z)$  in (6.3.11) can be obtained as

$$\Xi^*(z) = [\beta_{\perp} \bar{\beta}'_{\perp} \Theta^*(z) + \bar{\beta} \beta' B(z)^{-1}] \tag{6.3.15}$$

(see the proof of Proposition 6.1). The representation (6.3.11) will turn out to be useful, for example, in Chapter 9, where structural VECMs are discussed. The coefficient matrices  $\Xi_j^*$  of the operator  $\Xi^*(z)$  will then play an important role as specific impulse response coefficients. ■